# A System of Nonlinear Differential-Difference Equations with Variable Coefficients and Its Reductions

Sheng Zhang<sup>1, a</sup> and Xin Guo<sup>1,b</sup>

<sup>1</sup>School of Mathematics and Physics, Bohai University, Jinzhou 121013, PR China <sup>a</sup>szhangchina@126.com, <sup>b</sup>865959684@qq.com

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**Abstract.** Constructing integrable systems is a significant direction in soliton theory. In this paper, a new system of nonlinear differential-difference equations with variable coefficients is derived by introducing some derivable functions to the corresponding discrete spectral problems. In order to give some special cases of the derived differential-difference equations, three reductions are obtained which include Hirota's lattice equations as special cases. The processes of constructing such a system of variable-coefficient differential-difference equations and obtaining its reductions provide with a necessary help for the beginners.

## Introduction

Soliton equations are a kind of special nonlinear partial differential equations, one important property of such type of equations is the existence of so-called soliton solutions (stable waves) which spread in time without changing their size or shape and interact with each other in a particle-like way [1]. There is a close relation between the existence of soliton solutions and the integrability of equations, the known research results show that all the integrable systems exist soliton solutions [2, 3]. With the development of soliton theory, constructing integrable systems becomes a significant direction in nonlinear science [4-14]. What is integrable? there is still not a unified concept. Several integrabilities have been defined, such as Lax integrability, Liouville integrability and Painlevé integrability. Generally speaking, one should indicate the system is integrable in what sense.

Since the variable-coefficient systems could describe more realistic physical phenomena than their constant-coefficient counterparts when the inhomogeneities of media and nonuniformities of coundaries are taken into account [15, 16], we shall construct in the present paper a system of nonlinear differential-difference equations with variable coefficients.

## Derivation

Firstly, following the steps in [17] we consider the discrete spectral problems

$$\varphi_{1,n+1} = (1 - T_n S_n)^{-1} (z + R_n S_n) \varphi_{1,n} + (1 - T_n S_n)^{-1} (Q_n + \frac{1}{z} S_n) \varphi_{2,n},$$
(1)

$$\varphi_{2,n+1} = (1 - T_n S_n)^{-1} (R_n + z T_n) \varphi_{1,n} + (1 - T_n S_n)^{-1} (\frac{1}{z} + T_n Q_n) \varphi_{2,n}, \qquad (2)$$

with  $\varphi_{1,n}$  and  $\varphi_{2,n}$  satisfying the evolution equations

$$\frac{\partial}{\partial t}\varphi_{1,n} = A_n\varphi_{1,n} + B_n\varphi_{2,n}, \qquad (3)$$

$$\frac{\partial}{\partial t}\varphi_{2,n} = C_n\varphi_{1,n} + D_n\varphi_{2,n} \,. \tag{4}$$

Introducing shift operator *E* defined by  $E(\varphi_{i,n}) = \varphi_{i,n+1}$  and using the relation  $\frac{\partial}{\partial t}(E\varphi_{i,n}) = E(\frac{\partial \varphi_{i,n}}{\partial t})$  for i = 1, 2, we have

$$\frac{\partial}{\partial t}(E\varphi_{1,n}) = [(1 - T_n S_n)_t^{-1}(z + R_n S_n) + (1 - T_n S_n)^{-1}(z + R_n S_n)_t + (1 - T_n S_n)^{-1}(z + R_n S_n)A_n + (1 - T_n S_n)^{-1}(Q_n + \frac{1}{z}S_n)C_n]\varphi_{1,n} + [(1 - T_n S_n)^{-1}(z + R_n S_n)B_n + (1 - T_n S_n)_t^{-1}(Q_n + \frac{1}{z}S_n) + (1 - T_n S_n)^{-1}(Q_n + \frac{1}{z}R_n)_t + (1 - T_n S_n)^{-1}(Q_n + \frac{1}{z}S_n)D_n]\varphi_{2,n},$$
(5)

$$E(\frac{\partial}{\partial t}\varphi_{1,n}) = [(1 - T_n S_n)^{-1} (z + R_n S_n) A_{n+1} + (1 - T_n S_n)^{-1} (R_n + zT_n) B_{n+1}]\varphi_{1,n} + [(1 - T_n S_n)^{-1} (Q_n + \frac{1}{z} S_n) A_{n+1} + (1 - T_n S_n)^{-1} (\frac{1}{z} + Q_n T_n) B_{n+1}]\varphi_{2,n},$$
(6)

$$\frac{\partial}{\partial t}(E\varphi_{2,n}) = [(1 - T_n S_n)_t^{-1}(R_n + zT_n) + (1 - T_n S_n)^{-1}(R_n + zT_n)_t + (1 - T_n S_n)^{-1}(R_n + zT_n)A_n + (1 - T_n S_n)^{-1}(\frac{1}{z} + Q_n T_n)C_n]\varphi_{1,n} + [(1 - T_n S_n)_t^{-1}(\frac{1}{z} + Q_n T_n) + (1 - T_n S_n)^{-1}(Q_n T_n)_t + (1 - T_n S_n)^{-1}(Q_n + \frac{1}{z}R_n)_t + (1 - T_n S_n)^{-1}(Q_n + \frac{1}{z}S_n)D_n]\varphi_{2,n},$$
(7)

$$E(\frac{\partial}{\partial t}\varphi_{2,n}) = [(1 - T_n S_n)^{-1} (z + R_n S_n) C_{n+1} + (1 - T_n S_n)^{-1} (R_n + zT_n) D_{n+1}]\varphi_{1,n} + [(1 - T_n S_n)^{-1} (Q_n + \frac{1}{z} S_n) C_{n+1} + (1 - T_n S_n)^{-1} (\frac{1}{z} + Q_n T_n) D_{n+1}]\varphi_{2,n}.$$
(8)  
Eurther taking

Further taking

$$A_{n} = a_{n}^{(0)}(t) + za_{n}^{(1)}(t) , \quad B_{n} = b_{n}^{(0)}(t) + \frac{1}{z}b_{n}^{(-1)}(t) , \quad C_{n} = c_{n}^{(0)}(t) + xc_{n}^{(1)}(t) , \quad D_{n} = d_{n}^{(0)}(t) + \frac{1}{z}d_{n}^{(-1)}(t) , \quad (9)$$

then using Eq. (9) to collect the coefficients of z with same order in Eqs. (6-9) yields

$$z^{2}: \Delta_{n}a_{n}^{(1)}(t) = 0, \ a_{n}^{(1)}(t)T_{n} = c_{n+1}^{(1)}(t),$$
(10)

$$z^{-2}: d_n^{(-1)}(t)S_n = b_{n+1}^{(-1)}(t), \ \Delta_n d_n^{(-1)}(t) = 0,$$
(11)

$$z^{1}: b_{n}^{(0)}(t) = a_{n+1}^{(1)}(t)Q_{n}, \quad b_{n}^{(0)}(t)T_{n} = c_{n+1}^{(1)}(t)Q_{n},$$
(12)

$$z^{-1}: c_n^{(0)}(t)S_n = b_{n+1}^{(-1)}(t)R_n, \ c_n^{(0)}(t) = d_{n+1}^{(-1)}(t)R_n,$$
(13)

from which we have

$$a_n^{(1)}(t) = a(t), \ c_n^{(1)}(t) = a(t)T_{n-1}, \ b_n^{(-1)}(t) = d(t)S_{n-1},$$
(14)

$$d_n^{(-1)}(t) = d(t), \quad b_n^{(0)}(t) = a(t)Q_n, \quad c_n^{(0)}(t) = d(t)R_n.$$
(15)  
If we select

$$d_{n+1}^{(0)} = a^{(1)}(t)T_nQ_{n+1}, \ a_{n+1}^{(0)} = d^{(-1)}(t)S_nR_{n+1},$$
(16)

the following system of nonlinear differential-difference equations with variable coefficients can be obtained

$$R_{n,t} = (1 - R_n Q_n) [d(t)T_n - a(t)T_{n-1}],$$
(17)

$$S_{n,t} = (1 - S_n T_n) [a(t)Q_{n+1} - d(t)Q_n],$$
(18)

$$Q_{n,t} = (1 - R_n Q_n) [a(t)S_n - d(t)S_{n-1}],$$
(19)

$$T_{n,t} = (1 - S_n T_n) [d(t) R_{n+1} - d(t) R_n].$$
(20)

In particularly, when setting

$$R_n = 0, \ T_n = 1, \ Q_n = a(t)\beta_n, \ S_n = 1 - a(t)\alpha_n, \ a(t) = d(t),$$
(21)

Eqs. (17-20) can be reduced as a generalized Hirota's lattice equations with variable coefficients:

$$\beta_{n,t} = a(t)(\alpha_{n-1} - \alpha_n) - \frac{a'(t)}{a(t)}\beta_n, \qquad (22)$$

$$\alpha_{n,t} = a^2(t)\alpha_n(\beta_n - \beta_{n+1}) - \frac{a'(t)}{a(t)}\alpha_n.$$
(23)

which including the known Hirota's lattice equations [17]:

 $\beta_{n,t} = \alpha_{n-1} - \alpha_n, \tag{24}$ 

$$\alpha_{n,t} = \alpha_n (\beta_n - \beta_{n+1}) \,. \tag{25}$$

as special case as long as letting a(t) = 1.

If setting

$$\beta_n = -x_{n,t}, \ \alpha_n = e^{x_n - x_{n+1}}, \tag{26}$$

from Eqs. (22) and (23) we have

$$\beta_{n,t} = -x_{n,tt} = a(t)(e^{x_{n-1}-x_n} - e^{x_n-x_{n+1}}) + \frac{a'(t)}{a(t)}x_{n,t},$$
(27)

$$\alpha_{n,t} = (x_{n,t} - x_{n+1,t}) e^{x_n - x_{n+1}} = a^2(t) e^{x_n - x_{n+1}} (-x_{n,t} + x_{n+1,t}) - \frac{a'(t)}{a(t)} e^{x_n - x_{n+1}},$$
(28)

which can be reduced as

$$x_{n,tt} + a(t)(e^{x_{n-1}-x_n} - e^{x_n-x_{n+1}}) + \frac{a'(t)}{a(t)}x_{n,t} = 0,$$
(24)

$$x_{n,t} - x_{n+1,t} + \frac{a'(t)}{a^3(t) + a(t)} = 0.$$
<sup>(25)</sup>

Integrating Eq. (25) with respect to t one, we have

$$x_n - x_{n+1} = \log \frac{[1 + a^2(t)]^{\frac{1}{2}}}{a(t)},$$
(26)

then substituting Eq. (26) into Eq. (24) yields

$$x_{n,tt} + \frac{a'(t)}{a(t)} x_{n,t} = 0.$$
(27)

Solving Eq. (27), we obtain a special solution of Eqs. (24) and (25):

$$x_n = c_0 + c_1 \int \frac{1}{a(t)} dt \,.$$
(28)

## Conclusion

By introducing some derivable functions to the corresponding discrete spectral problems, a system of nonlinear differential-difference equations with variable coefficients is derive. To the best of our knowledge, this system of differential-difference equations has not been reported in literature. Besides, three reductions of the derived system of differential-difference equations are obtained which include Hirota's lattice equations as special cases. How to employ the method used in this paper to construct some other nonlinear differential-difference equations with variable coefficients is worthy of study. This is our tasks in the future.

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