

# Vibration Transfer Matrix Model of Multilayered Wave-attenuation Media of Warship

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**Abstract-Warship vibration isolation technology of multilayered wave-attenuation is based on new mechanism. In view of the influence of vibration transfer efficiency, which belongs to statics contains, this paper presents a two-dimensional transfer matrix model in the condition of oblique incidence, and gets derivation formulas of transmission and reflection coefficient. The construction of the transfer matrix model is based on the representation analysis of structure vibration energy and acoustic boundary conditions, and the new model is to study the implementation effect on vibration isolation technology of multilayered wave-attenuation. Then, by designing specialized experimental equipment and analyzing the results, the excellent implementation effect on multilayered wave-attenuation vibration isolation is verified. In addition, the calculation accuracy of the new model has been verified by contrastively analyzing the results of the new model with other literature.**

**Keywords-Multilayered Media; Transfer Matrix; Vibration Energy; Reflection Coefficient; Transmission Coefficient**

## I. INTRODUCTION

With continuous improvement of the status of marine rights and interests in national strategy, warship is to develop in the direction of large-scale, high-speed, integration, and concealment. Now underwater acoustic equipment operating distance is more and more far, the application of acoustic homing and fuzing weapons are becoming more and more widely, putting forward higher requirements for ship acoustic stealth technology. Therefore, navies bring the warship acoustic stealth technology into the focus of the research in high speed of arms race. At intermediate or low ship speed, main structural vibrations are originated from mechanical equipment, shafting system and weapon launcher motion, which are transmitted through supporting foundation, piping system and air to the hull, and then, the excited hull vibrations radiate wave energy into

water[1]. Research shows that, main energy of structural vibrations transmit through supporting foundation, and the most effective method to control it is to implement vibration isolation technology[2].

Vibration isolation technology has excellent characteristics of high damping capacity and convenient installation, which makes scholars focus the research on core actuators and vibration isolator development, and a lot of achievements have been obtained [3-4]. He et al [5] independently developed pneumatic vibration isolators for marine main engine, which had successfully realized the engineering application, and among them, type JYQN had excellent performance and applicability. Zhang et al [5] filled pneumatic vibration isolator with elastic and liquid to overcome problems of air leakage and instability, which lifted the restrictions of position monitoring and maintaining equipment. Because of filling vibration isolator with solid-liquid mixed media possessing obvious advantages in low-frequency and heavy-duty vibration isolation, a lot of literatures carry out researches on this field[7-9], and it is obvious that vibration isolation technology making full use of mixed media will be applied widely.

Literatures [10-11] put forward a new method to improve warship vibration isolation efficiency by filling closed isolator with multilayered wave-attenuation mixed media, thus, the vibration wave energy can be attenuated to the maximum extent. Comparing with filling single medium, the distinctive features of the new method act multiple wave-attenuation mechanisms comprehensively, including reflection, transmission, absorption, exciting liquid surface wave, etc. Therefore, researchers can say that the new method is able to achieve the purpose of high efficiency vibration isolation on root. Therefore, in order to provide important measure index and theoretical basis of wave-attenuation efficiency for improving vibration isolation theory,

Researchers present a two-dimensional transfer matrix and get derivation formulas of transmission and reflection coefficient.

## II. TRANSFER MATRIX MODEL DERIVATION

### A. Acoustic energy relationship and characterization

Medium volume element in sound field is  $\Delta V$ , sound pressure is  $P_0$ , density is  $\rho_0$ , volume element kinetic energy and potential Energy are respectively

$$E_k = \frac{1}{2} \Delta V \rho_0 v^2 \quad \text{and} \quad E_p = \frac{\Delta V}{2 \rho_0 c_0^2} P^2,$$

where  $v$  is particle vibration velocity,  $c_0$  is sound velocity in medium. Therefore, sound energy of volume element can be shown as,

$$E = E_k + E_p = \frac{\Delta V \rho_0}{2} (v^2 + \frac{P^2}{\rho_0^2 c_0^2}) \quad (1)$$

Sound energy density equation is,

$$\varepsilon = \frac{E}{\Delta V} = \frac{\rho_0}{2} (v^2 + \frac{P^2}{\rho_0^2 c_0^2}) \quad (2)$$

Let sound pressure expression be  $p(t, x) = p_a e^{j(\omega t - kx)}$ , particle vibration velocity be  $v(t, x) = v_a e^{j(\omega t - kx)}$ , among expressions,  $P_a$  and  $v_a$  are amplitudes of pressure expression and particle vibration velocity. Then, real parts of  $p(t, x)$  and  $v(t, x)$  are substituted into Formula 1, and instantaneous value of sound energy in unit volume is described as Formula 3,

$$\begin{aligned} E &= \frac{\Delta V \rho_0}{2} [\frac{P_a^2}{\rho_0^2 c_0^2} \cos^2(\omega t - kx) + \frac{P_a^2}{\rho_0^2 c_0^2} \cos^2(\omega t - kx)] \\ &= \frac{\Delta V P_a^2}{\rho_0 c_0^2} \cos^2(\omega t - kx) \end{aligned} \quad (3)$$

Average values of sound energy and sound energy

density during period  $T$  are

$$\bar{E} = \frac{1}{T} \int_0^T E dt = \frac{\Delta V P_a^2}{2 \rho_0 c_0^2},$$

$$\bar{\varepsilon} = \frac{P_a^2}{2 \rho_0 c_0^2}.$$

According to the definition of sound intensity, which is average acoustic energy flow of unit area in vertical direction, the relationship among sound intensity of sound energy, average value of sound energy and average value of sound energy density are shown as Formula 4,

$$I = \frac{c_0}{\Delta V T} \int_0^T E dt = \frac{c_0}{\Delta V} \bar{E} = c_0 \bar{\varepsilon} = \frac{P_a^2}{2 \rho_0 c_0} = \frac{\rho_0 c_0 v_a^2}{2} \quad (4)$$

In the application of practical engineering, total vibration energy at steady state and average value of sound energy density are generally used to evaluate structures vibration, but both of them cannot be measured directly. Thus, vibration acceleration is chosen

to be the measuring parameter in vibration level assessment. In the interval of  $(t_1, t_2)$ , mean square value

of energy signal is  $A = \int_{t_1}^{t_2} a^2(t) dt$ , and the damping effect of the vibration isolation measures can be described as Formula 5, which is defined as VLDA (vibration level difference of acceleration).

$$T = 10 \lg \frac{\frac{1}{n} \sum_{i=1}^n A_i^2}{\frac{1}{n} \sum_{i=1}^n A_{0i}^2} = 10 \lg \frac{\bar{A}^2}{\bar{A}_0^2} \quad (5)$$

Among it,  $\bar{A}^2$  is mean square value of post-isolation vibration acceleration,  $\bar{A}_0^2$  is mean square value of pre-isolation vibration acceleration. When single degree of freedoms structure is in forced vibration, the expression of vibration energy at steady state can be

described as  $W = \frac{m \bar{A}^2}{\omega^2}$ , among it,  $m$  is the mass of single degree of freedoms structure, and  $\omega$  is circular frequency of external force. Therefore, by modeling a two-dimensional transfer matrix in the condition of oblique incidence and getting transmission coefficient, the energy exchange relations of multilayered wave-attenuation mixed media can be obtained. It is obvious that sound intensity transmission coefficient is an important index to measure wave-attenuation efficiency.

### B. Modeling transfer matrix of multilayered media

Longitudinal wave can be generated in multilayered fluid system with plane wave oblique incidence, and acoustic boundary conditions, which continuities of sound pressure and normal mode velocity, are satisfied based on analysis of Fig .1. Therefore, the sound pressure field of layer  $i$  and layer  $i+1$  can be described as Formula 6 and Formula 7, among it,  $i$  represents the layer number, and it is to use  $D_{i-1}$  as coordinate axis  $x$  of layer  $i$ ,  $P^i = P_i^i + P_r^i$ ,  $P^{i+1} = P_i^{i+1} + P_r^{i+1}$ ,

$$\begin{cases} P_i^i = p_{ta}^i e^{j[\omega t - k_i(x - D_{i-1}) \cos \theta_i^i - k_i y \sin \theta_i^i]} \\ P_r^i = p_{ra}^i e^{j[\omega t + k_i(x - D_{i-1}) \cos \theta_r^i - k_i y \sin \theta_r^i]} \end{cases} \quad (6)$$

$$\begin{cases} P_i^{i+1} = p_{ta}^{i+1} e^{j[\omega t - k_{(i+1)}(x - D_i) \cos \theta_i^{i+1} - k_{(i+1)} y \sin \theta_i^{i+1}]} \\ P_r^{i+1} = p_{ra}^{i+1} e^{j[\omega t + k_{(i+1)}(x - D_i) \cos \theta_r^{i+1} - k_{(i+1)} y \sin \theta_r^{i+1}]} \end{cases} \quad (7)$$

According to Formula 6, Formula 7 and motion equations as Formula 8, the velocity potential can be obtained which are described as Formula 9 and Formula 10,

$$\begin{cases} v_x = -\frac{1}{\rho} \int \frac{\partial P}{\partial x} dt = \frac{\cos \alpha}{\rho c} P \\ v_y = -\frac{1}{\rho} \int \frac{\partial P}{\partial y} dt = \frac{\cos \beta}{\rho c} P \end{cases} \quad (8)$$

$$V^i = V_i^i + V_r^i = v_{ta}^i e^{j[\omega t - k_i(x - D_{i-1}) \cos \theta_i^i - k_i y \sin \theta_i^i]} + v_{ra}^i e^{j[\omega t + k_i(x - D_{i-1}) \cos \theta_i^i - k_i y \sin \theta_i^i]} \quad (9)$$

$$V^{i+1} = V_i^{i+1} + V_r^{i+1} = v_{ta}^{i+1} e^{j[\omega t - (k_{i+1} \cos \theta_{i+1}^i - k_i \cos \theta_i^i) x - (k_{i+1} \sin \theta_{i+1}^i - k_i \sin \theta_i^i) y]} + v_{ra}^{i+1} e^{j[\omega t + (k_{i+1} \cos \theta_{i+1}^i - k_i \cos \theta_i^i) x + (k_{i+1} \sin \theta_{i+1}^i - k_i \sin \theta_i^i) y]} \quad (10)$$

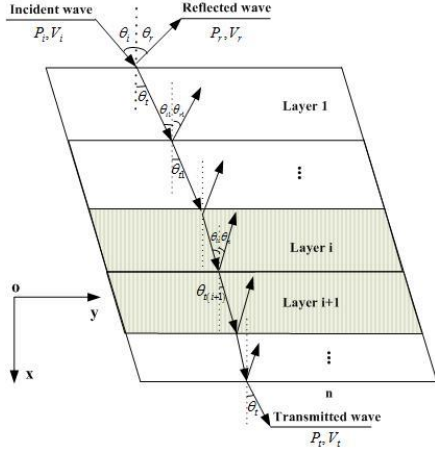


Figure 1. Sketch map of multilayered system subjected to oblique incidence

On the simultaneous of Formula 6 and Formula 7, Formula 9 and Formula 10, important equations can be obtained, which are  $Ap_{ta}^i + Bp_{ra}^i = p_{ta}^{i+1} + p_{ra}^{i+1}$  and  $Av_{ta}^i + Bv_{ra}^i = v_{ta}^{i+1} + v_{ra}^{i+1}$ , among them,  $A = e^{-jk_i d_i \cos \theta_i^i}$ ,  $B = e^{jk_i d_i \cos \theta_i^i}$ ,  $d_i = D_i - D_{i-1}$ . The relations of sound pressure and particle vibration velocity in layer  $i$  and layer  $i+1$  is as follows,

$$\begin{cases} p_{ta}^{i+1} = A \left( \frac{Z_{i+1}}{Z_i} + 1 \right) p_{ta}^i + B \left( 1 - \frac{Z_{i+1}}{Z_i} \right) p_{ra}^i \\ p_{ra}^{i+1} = A \left( 1 - \frac{Z_{i+1}}{Z_i} \right) p_{ta}^i + B \left( \frac{Z_{i+1}}{Z_i} + 1 \right) p_{ra}^i \end{cases} \quad (11)$$

$$\begin{cases} v_{ta}^{i+1} = A \left( \frac{Z_i}{Z_{i+1}} + 1 \right) v_{ta}^i + B \left( \frac{Z_i}{Z_{i+1}} - 1 \right) v_{ra}^i \\ v_{ra}^{i+1} = A \left( \frac{Z_i}{Z_{i+1}} - 1 \right) v_{ta}^i + B \left( \frac{Z_i}{Z_{i+1}} + 1 \right) v_{ra}^i \end{cases} \quad (12)$$

$$Z_i = \frac{Ri}{\cos \theta_i^i} = \frac{\rho_i c_i}{\cos \theta_i^i}$$

Among them, the characteristic impedance of layer  $i$ ,  $\theta_i^i = \theta_r^i$ ,  $\theta_i^{i+1} = \theta_r^{i+1}$ ,  $\sin \theta_i^i / \sin \theta_i^{i+1} = k_{i+1} / k_i = c_i / c_{i+1}$ ,  $k_i$  is the wave number of layer  $i$ . Modifying Formula 11 and Formula

$$\begin{bmatrix} p_{ta}^{i+1} \\ v_{ta}^{i+1} \end{bmatrix} = \frac{1}{2} \mathbf{H} \begin{bmatrix} p_{ta}^i \\ v_{ta}^i \end{bmatrix},$$

$$\begin{bmatrix} p_{ra}^{i+1} \\ v_{ra}^{i+1} \end{bmatrix} = \frac{1}{2} \mathbf{J} \begin{bmatrix} p_{ra}^i \\ v_{ra}^i \end{bmatrix},$$

$\mathbf{H}$  and  $\mathbf{J}$  are transfer matrix of transmission and reflection coefficient in layer  $i$  and layer  $i+1$ , to expand  $\mathbf{H}$  and  $\mathbf{J}$  as follows:

$$\mathbf{H} = \begin{bmatrix} \left( \frac{Z_{i+1}}{Z_i} + 1 \right) e^{-jk_i d_i \cos \theta_i^i} & \left( 1 - \frac{Z_{i+1}}{Z_i} \right) e^{jk_i d_i \cos \theta_i^i} \\ \left( \frac{Z_i}{Z_{i+1}} + 1 \right) e^{-jk_i d_i \cos \theta_i^i} & \left( \frac{Z_i}{Z_{i+1}} - 1 \right) e^{jk_i d_i \cos \theta_i^i} \end{bmatrix} \quad (13)$$

$$\mathbf{J} = \begin{bmatrix} \left( 1 - \frac{Z_{i+1}}{Z_i} \right) e^{-jk_i d_i \cos \theta_i^i} & \left( 1 + \frac{Z_{i+1}}{Z_i} \right) e^{jk_i d_i \cos \theta_i^i} \\ \left( \frac{Z_i}{Z_{i+1}} - 1 \right) e^{-jk_i d_i \cos \theta_i^i} & \left( \frac{Z_i}{Z_{i+1}} + 1 \right) e^{jk_i d_i \cos \theta_i^i} \end{bmatrix} \quad (14)$$

### III. SOUND INTENSITY TRANSFER COEFFICIENT

In the multilayered system subjected to normal incidence plane waves, which is  $\theta_i = 0$ ,

transfer matrix results of Formula 13 and Formula 14 in layer  $i$  and layer  $i+1$  have the same form with sound pressure transfer matrix  $A_i$  of literature [12]. Then, rewriting

$$\mathbf{H}^i = \begin{bmatrix} H_{11}^i & H_{12}^i \\ H_{21}^i & H_{22}^i \end{bmatrix},$$

Formula 13 into a general form as transfer function of  $n$  layers can be described as,

$$p_{ra}^n = \left( \frac{1}{2} \right)^n H_{21}^1 \cdot H_{21}^2 \cdots H_{21}^n p_{ta}^0 + \left( \frac{1}{2} \right)^n H_{22}^1 \cdot H_{22}^2 \cdots H_{22}^n p_{ra}^0 \quad (15)$$

Then, substituting the known conditions as  $p_{ra}^0 = P_r$ ,  $p_{ta}^0 = P_i$ ,  $p_{ra}^n = 0$  into Formula 14, sound pressure reflection coefficient  $\xi$  and transmission coefficient  $\eta$  can be resolve to,

$$\xi = \frac{|P_r|^2 / 2R_1}{|P_i|^2 / 2R_1} = \left| \frac{\prod_{l=1}^n H_{21}^l}{\prod_{l=1}^n H_{22}^l} \right|^2 \quad (16)$$

$$\eta = \frac{|P_t|^2 / 2R_n}{|P_i|^2 / 2R_1} = \left( \frac{1}{2} \right)^{2n} \frac{R_1}{R_n} \left| \prod_{l=1}^n H_{11}^l - \frac{\prod_{l=1}^n (H_{12}^l \cdot H_{21}^l)}{\prod_{l=1}^n H_{22}^l} \right|^2 \quad (17)$$

In the process of sound waves transfer in fluid, the medium features of viscosity, heat conduction and relaxation effect show significant energy attenuation, thus, it is needed to turn the expressions of sound velocity and wave number which are closely related to characteristic impedance into form of complex.

Complex sound velocity is  $\tilde{c} = c(1 - i\varepsilon_c)$  and complex wave number is  $\tilde{k} = k(1 + i\varepsilon_k)$ , among them,  $\varepsilon_c$  and  $\varepsilon_k$  are the loss factors.

#### IV. EXPERIMENTAL RESULTS AND MODEL TEST

According to derivations and actual measuring demand, the test system of vibration isolation technology of multilayered wave-attenuation is designed, the schematic diagram of which is shown in Fig .2. The exciter is precision speed regulating motor of LN30/2 with vibration frequency of 0-300Hz, and the closed isolator is filled with multilayered wave-attenuation media of water and dimethicone. Upper and lower mass both are 5Kg rigid structures, vibration sensor adopts three-axis acceleration chip with adjustable measuring range, and the testing equipment is the STC12C5A60S2 series MCU. Fig .3 shows the VLDA of filling isolator with different media. From Fig .3, it can be seen that the VLDA increases firstly and then decreases with the increase of water filling amount, when water mass reaches 0.4kg, the max VLDA is 9.61dB. The VLDA variation filling with dimethicone is consistent with that of filling with water, and the max VLDA is 10.2dB. When filling with mixed media of water and dimethicone, the max VLDA reaches 14.08dB at 0.5kg total mass, thus, filling with multilayered wave-attenuation media has the excellent implementation effect on vibration isolation.

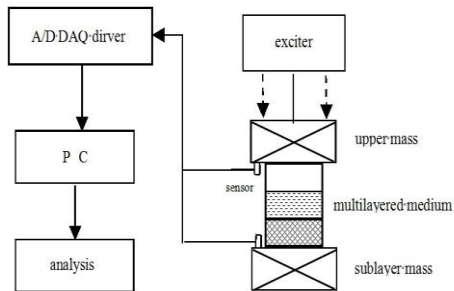


Figure 2. The equipment of test system

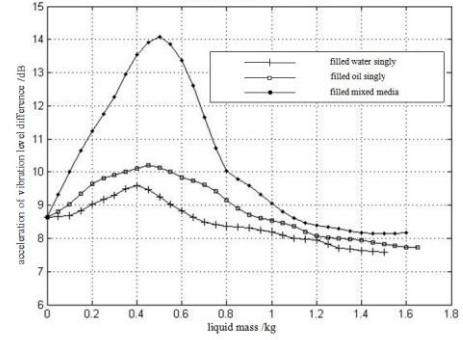


Figure 3. Vibration level difference of acceleration at different media

Using the parameter data shown in Tab 1, absorption coefficient of literature [13] and two-dimensional transfer matrix model in wave-attenuation media can be resolved, then, showing the comparison results in Fig .4. As can be seen from Fig .4, the calculation results of new model are consistent with that of literature [13]. Thus, the rationality of two-dimensional transfer matrix model is verified.

TABLE I. PARAMETER TABLE

| Position  | 1 <sup>st</sup> Layer | 2 <sup>nd</sup> Layer | 3 <sup>rd</sup> Layer           |
|---|-----------------------|-----------------------|---------------------------------|
| Materials   | Dimethicone           | Water                 | CH <sub>2</sub> Cl <sub>2</sub> |
| Temperature<br>°C   | 20                    | 20                    | 20                              |
| Sound Velocity<br>×10 <sup>3</sup> m/s                          | 1.040                 | 1.483                 | 1.070                           |
| Density<br>×10 <sup>3</sup> kg/m <sup>3</sup>                   | 0.963                 | 0.998                 | 1.330                           |
| Characteristic Impedance<br>×10 <sup>6</sup> N·s/m <sup>3</sup> | 1.001                 | 1.480                 | 1.423                           |
| loss factor<br>×10 <sup>-2</sup>                                | 2                     | 0.1                   | 5                               |

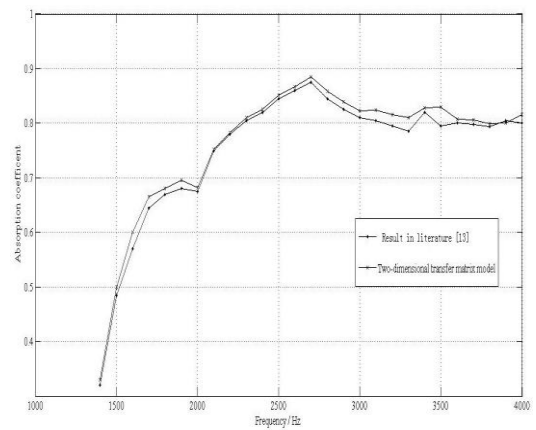


Figure 4. Comparison figure of calculation results

#### V. CONCLUSIONS

In this paper, the relationship of vibration energy and acoustic field are analyzed, and the conclusion that sound intensity transfer coefficient can be used as the standard for vibration attenuation evaluation of multilayered media is obtained, and then, general

formulas of transfer coefficient are deduced based on the representation analysis of structure vibration energy and acoustic boundary conditions; The excellent implementation effect on multilayered wave-attenuation vibration isolation is verified by specialized experimental research; Meanwhile, the calculation accuracy of the new model has been verified by comparative analyzing the results of the new model with other literature. In a word, the study not only creates favorable conditions for improving vibration isolation efficiency, but also provides a theoretical and experimental basis for improving vibration isolation theory and developing new isolators.

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