

New Exact Periodic Wave Solutions of the Nonlinear Klein-gordon Equation

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Abstract -To the non-linear Klein-Gordon equations, the undetermined assumption method is used to get the exact periodic wave solution in the form of Jacobi elliptic function fraction and with the asymptotic values non-zero. The condition for their existence and boundness are obtained and the impact of the change of travelling periodic wave velocity is revealed upon the periodic wave solution and the size of the period.

Keywords-Nonlinear Klein-Gordon equation, Jacobi elliptic function, Periodic wave solutions

I. INTRODUCTION

Non-linear Klein-Gordon equation^[1]

$$u_{tt} - u_{xx} + \lambda u + \mu u^3 = 0 \quad (1)$$

is an important model in mechanics and particle physics. In [3,4], by expansion to the Jacobi elliptic function, the exact periodic solution is got to the equation in the following form.

$$u(\xi) = \sum_{i=0}^n a_i sn^i(\xi) + \sum_{j=1}^n b_j sn^{-j}(\xi)$$

Here, we try to give the periodic wave solution to the non-linear Klein-Gordon equation (1) with the form as

$$x(\xi) = \left[\frac{A_0 + A_1 cn^2(\alpha\xi, k)}{B_0 + B_1 cn^2(\alpha\xi, k)} \right]^{\frac{1}{2}} \quad (2)$$

and get to know the existence and bound condition of the solution.

II. NEW EXACT PERIODIC WAVE SOLUTION TO THE NON-LINEAR KLEIN-GORDON EQUATION

It is easily known that the travelling wave solution of equation (1) $u(x, t) = u(\xi) = u(x - vt)$ satisfies,

$$(v^2 - 1)u''(\xi) + \lambda u + \mu u^3 = 0$$

$$\text{or} \quad u''(\xi) + lu + su^3 = 0 \quad (3)$$

in which, $l = \frac{\lambda}{v^2 - 1}$, $s = \frac{\mu}{v^2 - 1}$

Let $u(\xi) = \sqrt{p(\xi)}$, then equation (3) can be changed as

$$\frac{1}{2} p(\xi) p''(\xi) - \frac{1}{4} (p'(\xi))^2 + lp^2(\xi) + sp^3(\xi) = 0. \quad (4)$$

To get the solution to equation (1), we just need the solution to equation (3), which is the positive solution to equation (4). Let equation (4) have the solution with the form as

$$p(\xi) = \frac{A_0 + A_1 cn^2(\alpha\xi, k)}{B_0 + B_1 cn^2(\alpha\xi, k)} \quad (\text{k is module}) \quad (5)$$

From (5), we can easily get the expression of $p'(\xi)$, $p''(\xi)$, and then put the expressions and (5) into (4). After simplification, and by comparing the coefficients of the powers of $cn(\alpha\xi, k)$ and setting them as zero, we get the following equations,

$$\begin{cases} A_0 B_0 (k^2 - 1)(A_0 B_1 - A_1 B_0) \alpha^2 + A_0^2 (A_0 B_0 s + B_0^2 l) = 0 \\ 2A_0 (2B_0 k^2 - B_0 + B_1 k^2 - B_1)(A_1 B_0 - A_0 B_1) \alpha^2 \\ \quad + 2A_0 B_0 (A_0 B_1 + A_1 B_0) l + A_0^2 (A_0 B_1 + 3A_1 B_0) s = 0 \\ (A_0 B_1 - A_1 B_0)((2A_0 B_1 - 3A_1 B_1 - 2A_1 B_0 + 3A_0 B_0) k^2 - A_0 B_1 \\ \quad + 3A_1 B_1 + A_1 B_0) \alpha^2 + (A_0^2 B_1^2 + 4A_0 A_1 B_0 B_1 + A_1^2 B_0^2) l \\ \quad + 3A_0 A_1 (A_0 B_1 + A_1 B_0) s = 0 \\ 2A_1 (B_0 k^2 + 2B_1 k^2 - B_1)(A_0 B_1 - A_1 B_0) \alpha^2 \\ \quad + 2A_1 B_1 (A_1 B_0 + A_0 B_1) l + A_1^2 (A_1 B_0 + 3A_0 B_1) s = 0 \\ A_1 B_1 k^2 (A_0 B_1 - A_1 B_0) \alpha^2 - A_1^2 (A_1 B_1 s + B_1^2 l) = 0 \end{cases} \quad (6)$$

To solve the equations with the Maple, we get

$$A_0 = 0, B_0 = \frac{2y(1 - 2k^2)(ly + s)}{(4ly + 3s)k^2} A_1, B_1 = yA_1, \quad (1)$$

$$\alpha^2 = \frac{4ly + 3s}{2y(2k^2 - 1)}. \quad (7)$$

Here y satisfies the equation

$$4l^2 y^4 + 6lsy^3 - s^2 (k^4 - k^2 - 2)y^2 = 0. \quad (8)$$

Then we get the root of equation (8)

$$y_1 = -\frac{s(k^2 + 1)}{2l}, y_2 = \frac{s(k^2 - 2)}{2l}, y_{3,4} = 0.$$

Put y_i ($i = 1, 2, 3, 4$) respectively into (7), we get the expressions of A_i, B_i ($i = 0, 1$) and α , and then put them into (5), the solution to equation (4) is obtained. And put it into $u(\xi) = \sqrt{p(\xi)}$, we get the solution of equation (3), and further we get the solution of equation (1).

When $y = y_1$, there exists

$$u_1(\xi) = \left[\frac{2k^2 \lambda}{\mu(1+k^2)} \cdot \frac{cn^2\left(\sqrt{\frac{\lambda}{(v^2-1)(1+k^2)}}\xi, k\right)}{1-k^2+k^2cn^2\left(\sqrt{\frac{\lambda}{(v^2-1)(1+k^2)}}\xi, k\right)} \right]^{\frac{1}{2}} \quad (9)$$

Note: In the computation, just omit the meaningless solutions and the solutions determined as boundless according to the nature of Jacobi ellipse function or with $p(\xi)$ less than zero. We do it the way in the following and will not repeat this.

2)

$$A_0 = -A_1, B_0 = \frac{y(k^2 - 1)(2k^2ly - 2ly + k^2s - 2s)}{s(1+k^2)k^2},$$

$$B_1 = yA_1, \alpha^2 = -\frac{s^2(k^2 + 1)}{2y(k^2 - 1)^2(4ly + 3s)}.$$

in which y satisfies the equation

$$4l^2(k^2 - 1)^2y^4 + 6ls(k^2 - 1)^2y^3 + (2k^4 - 5k^2 + 2)s^2y^2 = 0. \quad (10)$$

Then we easily get the root of equation (10)

$$y_1 = \frac{(2 - k^2)s}{2(k^2 - 1)l}, y_2 = \frac{(1 - 2k^2)s}{2(k^2 - 1)l}, y_{3,4} = 0.$$

As the above mentioned, here when $y = y_2$, there exists

$u_2(\xi) =$

$$\left[\frac{2k^2 \lambda(1 - k^2)}{\mu(1 - 2k^2)} \cdot \frac{1 - cn^2\left(\sqrt{\frac{\lambda}{(1 - 2k^2)(v^2 - 1)}}\xi, k\right)}{1 - k^2 + k^2cn^2\left(\sqrt{\frac{\lambda}{(1 - 2k^2)(v^2 - 1)}}\xi, k\right)} \right]^{\frac{1}{2}} \quad (11)$$

$$A_0 = \frac{2l^2 - 2(k^4 - k^2 + 1)y^2}{3k^2sy} B_1,$$

3)

$$A_1 = 0, B_0 = -\frac{2k^2y - y + l}{3k^2y} B_1, \alpha^2 = y.$$

Here y satisfies the following equation

$$(k^2 - 2)(2k^2 - 1)(k^2 + 1)y^3 - 3l(k^4 - k^2 + 1)y^2 + l^3 = 0. \quad (12)$$

We can easily get the root of equation (12),

$$y_1 = \frac{l}{k^2 + 1}, y_2 = \frac{l}{1 - 2k^2}, y_3 = \frac{l}{k^2 - 2}.$$

And when $y = y_3$, there exists

$$u_3(\xi) = \left[\frac{2(k^2 - 1)\lambda}{\mu(2 - k^2)} \cdot \frac{1}{1 - k^2 + k^2cn^2\left(\sqrt{\frac{\lambda}{(k^2 - 2)(v^2 - 1)}}\xi, k\right)} \right]^{\frac{1}{2}} \quad (13)$$

$$A_0 = 0, A_1 = \frac{2lk^2}{s(1 - 2k^2)} B_0,$$

4)

$$B_1 = 0, \alpha^2 = \frac{l}{1 - 2k^2}.$$

With the same method, we can directly get

$$u_4(\xi) = \left[\frac{2k^2 \lambda}{\mu(1 - 2k^2)} \cdot cn^2\left(\sqrt{\frac{\lambda}{(1 - 2k^2)(v^2 - 1)}}\xi, k\right) \right]^{\frac{1}{2}}. \quad (14)$$

$$A_0 = -\frac{2l(1 - k^2)}{s(2 - k^2)} B_0,$$

$$5) A_1 = -\frac{2lk^2}{s(2 - k^2)} B_0,$$

$$B_1 = 0, \alpha^2 = -\frac{l}{2 - k^2}.$$

The same as the above, we can get

$$u_5(\xi) = \left[\frac{2\lambda(1 - k^2 + k^2cn^2\left(\sqrt{-\frac{\lambda}{(2 - k^2)(v^2 - 1)}}\xi, k\right))}{\mu(k^2 - 2)} \right]^{\frac{1}{2}} \quad (15)$$

$$A_0 = -\frac{2lk^2}{s(1+k^2)}B_0,$$

$$6) A_1 = \frac{2lk^2}{s(1+k^2)}B_0,$$

$$B_1 = 0, \alpha^2 = \frac{l}{1+k^2}.$$

Also we can get

$$u_6(\xi) =$$

$$\left[-\frac{2k^2\lambda}{\mu(1+k^2)} \cdot (1-cn^2 \left(\sqrt{\frac{\lambda}{(1+k^2)(v^2-1)}}\xi, k\right))\right]^{\frac{1}{2}} \quad (16)$$

Based on analysis and discussion, we know: $u_1(\xi)$ 、 $u_6(\xi)$ need to satisfies the condition $\lambda(v^2-1) > 0, \mu(v^2-1) < 0$, and $u_2(\xi)$ 、 $u_4(\xi)$ need to satisfies the condition $\lambda(1-2k^2)(v^2-1) > 0, \mu(v^2-1) > 0$ while $u_3(\xi)$ 、 $u_5(\xi)$ need the condition $\lambda(v^2-1) < 0, \mu(v^2-1) > 0$.

III. DISCUSSION ON THE PERIODIC WAVE SOLUTION TO EQUATIONS WITH THE FORM AS (2)

For convenience to discussion, we take it for example that equation (1) satisfies conditions

$$\lambda(v^2-1) > 0, \mu(v^2-1) < 0$$

which are the conditions that $u_1(\xi)$ 、 $u_6(\xi)$ satisfy.

When $\lambda > 0, \mu < 0$, obviously $v^2 > 1$, then equation

(1) has the large-velocity periodic wave solution $u_1(\xi)$ 、 $u_6(\xi)$ while, when $\lambda < 0, \mu > 0$, obviously $v^2 < 1$, the equation has small-velocity periodic wave solution $u_1(\xi)$ 、 $u_6(\xi)$. Analysis shows that the waveform and cycle of the periodic wave solution are influenced by its velocity v and that the bigger the velocity of the periodic wave solution, the bigger the period becomes with the change in the waveform.

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